

A Galerkin method for Uncertain Hyperbolic Systems: Roe solver and Entropy corrector.

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Workshop Torino
May 10-13, 2010



Introduction

Parametric Uncertainty Quantification

- uncertainties in **input quantities** (model parameters, initial and boundary conditions)
- uncertain quantities **parametrized** by random variables with known distribution functions

Stochastic spectral methods

- decompose random quantities on **suitable approximation bases** (Ghanem and Spanos 91)
- **Stochastic expansion of the solution of a model :**

$$U(x, t, \xi) \approx U^P(x, t, \xi) = \sum_{\alpha=1}^P u_{\alpha}(x, t) \Psi_{\alpha}(\xi).$$

$u_{\alpha}(x, t)$ **stochastic modes** of the solution.

Galerkin
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Stochastic hyperbolic
systems

Stochastic
discretization

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Galerkin system

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Numerical scheme

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The upwind scheme

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Periodic Burgers
equation

Euler equations

Entropy corrector

Different computational strategies

- **probabilistic collocation** : stochastic modes evaluated by polynomial interpolation
- **non-intrusive projection** : stochastic modes evaluated by numerical integration
- **stochastic Galerkin** : reformulated deterministic problem for the stochastic modes

Stochastic Galerkin methods :

- rely on a **weak form** of the problem
- well suited for **mathematical analysis**
- design of **adaptive methods**

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State of the art

Stochastic spectral methods applied to a large variety of engineering problems (elasticity, thermal science, fluid flows, chemical/biological systems,...) governed by elliptic, parabolic ,ODE or incompressible NS models.

Hyperbolic models

Two main difficulties :

- solutions **discontinuous** in the spatial as well as in the stochastic domains
- **nonlinearities** in the flux functions

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Hyperbolic models

- **non-intrusive approaches** : multi-element probabilistic collocation methods (Lin et al. 08)
- **pseudo-intrusive methods** : stochastic modes of flux computed by quadrature methods (Ge et al. 08, Poette et al. 09)
- **intrusive methods** : scalar linear wave equation (Gottlieb and Xiu 08), mean flux upwinding (Lin et al. 06)

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Objectives

- Develop an **intrusive stochastic Galerkin method**
- Investigate **hyperbolicity of the Galerkin system**
- Design a **Roe-type solver with entropy corrector**

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 - Hyperbolicity of the Galerkin system
- 2 Numerical method**
- 3 Results**

Stochastic parametrization

$\xi = (\xi_1, \dots, \xi_N) \sim \mathcal{U}(\Xi = [0, 1]^N) \rightarrow L^2(\Xi, p_\xi = 1)$
 corresponding space of the second-order random variables
 with the expectation operator $\langle H \rangle = \int_{\Xi} H(y) p_\xi(y) dy$.

Stochastic hyperbolic systems

We seek for $U(x, t, \xi) \in \mathbb{R}^m \otimes L^2(\Xi, p_\xi)$ solving a.s.

$$\begin{cases} \frac{\partial}{\partial t} U(x, t, \xi) + \frac{\partial}{\partial x} F(U(x, t, \xi); \xi) = 0, \\ U(t = 0, x, \xi) = U_0(x, \xi). \end{cases}$$

$$(x, t, \xi) \in \Omega \times [0, T] \times \Xi, \quad \Omega = [0, 1],$$

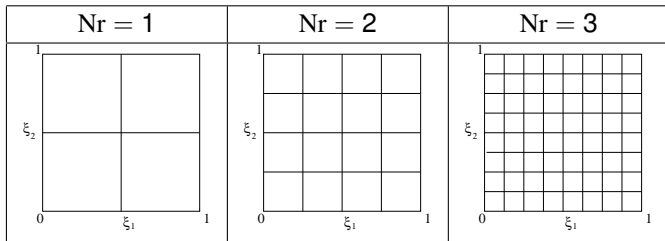
$\nabla_U F \in \mathbb{R}^{m,m} \otimes L^2(\Xi, p_\xi)$ stochastic Jacobian matrix
 \mathbb{R} -diagonalizable almost surely.

Stochastic discretization

We approximate $U(x, t, \xi)$ in the **stochastic space** of fully tensorized piecewise polynomial functions $\mathcal{S}^{\text{No}, \text{Nr}}$:

- Nr : resolution level $\rightarrow 2^{\text{NNr}}$ **SE** (stochastic elements),
- No : expansion order $\rightarrow U(\xi)|_{\text{SE}} \approx U^{\text{P}}(\xi)|_{\text{SE}} \in \mathbb{Q}_{\text{No}}^{\text{N}}[\xi]$.

Case $\text{N} = 2$.



$$\dim \mathcal{S}^{\text{No}, \text{Nr}} = (\text{No} + 1)^{\text{N}} 2^{\text{NNr}} := \text{P}.$$

Stochastic discretization

We approximate $U(x, t, \xi)$ in the stochastic space of fully tensorized piecewise polynomial functions $\mathcal{S}^{\text{No}, \text{Nr}}$:

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$$\dim \mathcal{S}^{\text{No}, \text{Nr}} = (\text{No} + 1)^{\text{Nr}} 2^{\text{Nr}} := \text{P}.$$

Select the Stochastic Element (SE) orthonormal basis $\{\Psi_{\alpha}\}_{\alpha=1, \dots, \text{P}}$, which corresponds to local fully tensorized (rescaled) Legendre polynomial bases, s.t.

$$\text{span}(\Psi_1, \dots, \Psi_{\text{P}}) = \mathcal{S}^{\text{No}, \text{Nr}}.$$

The Galerkin system

Stochastic expansion of the solution :

$$U(x, t, \xi) \approx U^P(x, t, \xi) = \sum_{\alpha=1}^P u_{\alpha}(x, t) \psi_{\alpha}(\xi).$$

$u_{\alpha}(x, t) \in \mathbb{R}^m$ stochastic modes of the solution.

Galerkin projection of the original stochastic problem :

$$\begin{cases} \left\langle \psi_{\alpha} \frac{\partial U^P}{\partial t} \right\rangle + \left\langle \psi_{\alpha} \frac{\partial F(U^P; \cdot)}{\partial x} \right\rangle = 0, & \forall \alpha = 1, \dots, P, \\ \langle \psi_{\alpha} U^P \rangle (t=0) = \langle \psi_{\alpha} U_0(x, \cdot) \rangle, & \forall \alpha = 1, \dots, P. \end{cases}$$

The Galerkin system

We seek for $u(x, t) \in \mathbb{R}^{mP}$ solving

$$\frac{\partial}{\partial t} u(x, t) + \frac{\partial}{\partial x} f(u(x, t)) = 0, \quad u(x, t = 0) = u^0(x),$$

$$u(x, t) = (u_1(x, t), \dots, u_P(x, t))^T \in \mathbb{R}^{mP},$$

$$f(u(x, t)) = (f_1(u), \dots, f_P(u))^T \in \mathbb{R}^{mP},$$

$$f_\alpha(u) := \langle \Psi_\alpha F(U^P; \cdot) \rangle, \quad \alpha = 1, \dots, P.$$

Is the Galerkin system hyperbolic ?

Hyperbolicity of the Galerkin system

$$(\nabla_u f(u))_{\alpha, \beta=1, \dots, P} = \langle \nabla_U F(U^P; \cdot) \Psi_\alpha \Psi_\beta \rangle_{\alpha, \beta=1, \dots, P}$$

Is $\nabla_u f(u) \in \mathbb{R}^{mP, mP}$ \mathbb{R} -diagonalizable ?

Diagonal block structure owing to the decoupling of the problem over different stochastic elements :

$$\nabla_u f(u) = \begin{pmatrix} [\nabla_u f]^{SE=1} & \dots & 0 & \dots & 0 \\ \vdots & \ddots & & & \vdots \\ 0 & \dots & [\nabla_u f]^{SE=i} & \dots & 0 \\ \vdots & & & \ddots & \vdots \\ 0 & \dots & 0 & \dots & [\nabla_u f]^{SE=2^{NNr}} \end{pmatrix}$$

→ It suffices to analyze one single block $[\nabla_u f]^{SE}$ of size $m(\text{No} + 1)^N$.

Hyperbolicity of the Galerkin system

HYPERBOLICITY proven in **two specific cases** :

Theorem

HYPERBOLICITY, if the original stochastic system is a *scalar conservation law*.

Theorem

HYPERBOLICITY, if the stochastic Jacobian matrix

$$\nabla_U F(\cdot; \xi)$$

- is either *symmetric* (almost surely)
- or its *eigenvectors* are *deterministic* (independent of the uncertainty).

Applications : Scalar wave equation with uncertain sound velocity. / Linear hyperbolic systems with uncertainty only on initial or boundary conditions.

Hyperbolicity of the Galerkin system

In the **general case**, we consider the **approximate Galerkin Jacobian matrix** $\left[\overline{\nabla_u f}\right]^{SE}$ whose coefficients are obtained by approaching the coefficients of $[\nabla_u f]^{SE}$ by a **Gauss quadrature**

$$\left[\left(\overline{\nabla_u f(u)} \right)_{\alpha, \beta=1, \dots, (No+1)^N} \right]^{SE} = \left(\sum_{\gamma=1}^{(No+1)^N} \varpi_{\gamma} \nabla_u F(U^P(\xi_{\gamma}); \xi_{\gamma}) \Psi_{\alpha}(\xi_{\gamma}) \Psi_{\beta}(\xi_{\gamma}) \right)_{\alpha, \beta=1, \dots, (No+1)^N},$$

$\{\xi_{\gamma}\}_{\gamma=1, \dots, (No+1)^N}$ set of the **Gauss points** in the SE, with associated weights $\{\varpi_{\gamma}\}_{\gamma=1, \dots, (No+1)^N}$.

Hyperbolicity of the Galerkin system

Assume $\nabla_U F(U^P(\xi); \xi) = L(\xi)\Lambda(\xi)R(\xi)$,

Theorem

$[\nabla_U f]^{SE}$ is \mathbb{R} -diagonalizable.

Explicit expressions of the eigenvalues

$$\{\lambda'_\gamma\} = \{\Lambda^k(\xi_\eta)\},$$

and of the left and right eigenvectors

$$\begin{aligned} \{r'_\gamma\} &= \{(\varpi_\eta R^k(\xi_\eta) \Psi_\beta(\xi_\eta))\}, \\ \{l'_\gamma\} &= \{(\varpi_\eta L^k(\xi_\eta) \Psi_\beta(\xi_\eta))\}. \end{aligned}$$

$\rightarrow \{\lambda'_\gamma\}$, $\{r'_\gamma\}$, and $\{l'_\gamma\}$, $\gamma = 1, \dots, m(\text{No} + 1)^N$,

approximations of the eigenvalues and eigenvectors of $[\nabla_U f]^{SE}$.

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Discretization of the Galerkin system using a **FV method** :

$$u_i^{n+1} = u_i^n - \frac{\Delta^n t}{\Delta x} (\varphi(u_i^n, u_{i+1}^n) - \varphi(u_{i-1}^n, u_i^n)),$$

- Δx (uniform) spatial step,
- $\Delta^n t$ time step,
- $u_i^n \approx \int_{(i-1/2)\Delta x}^{(i+1/2)\Delta x} u(x, t^n) dx,$
- $\varphi(\cdot, \cdot)$ 1st order numerical flux function :

$$\varphi(u_i^n, u_{i+1}^n) = \underbrace{\frac{f(u_i^n) + f(u_{i+1}^n)}{2}}_{\substack{\uparrow \\ \text{centered part of the flux}}} - \underbrace{|a(u_i^n, u_{i+1}^n)|}_{\substack{\uparrow \\ \text{upwind matrix}}} \frac{u_{i+1}^n - u_i^n}{2}.$$

Roe linearized matrix

- Assume that the original stochastic problem possesses a **Roe linearized matrix** and a **Roe state** a.s.,

$$(U_L, U_R) \rightarrow A^{\text{Roe}}(U_L, U_R) = \nabla_U F(U_{LR}^{\text{Roe}}; \cdot).$$

- Given two states u_L and u_R of the Galerkin system,

$$\left. \begin{array}{l} u_L \rightarrow U_L^P \\ u_R \rightarrow U_R^P \end{array} \right\} \rightarrow U_{LR}^{\text{Roe}},$$

$$\rightarrow a^{\text{Roe}}(u_L, u_R) = \langle \nabla_U F(U_{LR}^{\text{Roe}}; \cdot) \Psi_\alpha \Psi_\beta \rangle.$$

Theorem

a^{Roe} is a **Roe linearized matrix** for the Galerkin system.

Choice of upwinding :

$$\varphi(u_i^n, u_{i+1}^n) = \frac{f(u_i^n) + f(u_{i+1}^n)}{2} - |a^{\text{Roe}}(u_i^n, u_{i+1}^n)| \frac{u_{i+1}^n - u_i^n}{2}.$$

- **Consistency** of the numerical scheme
- **Conservativity** through shocks

Efficient approximation of $|a^{\text{Roe}}(u_i^n, u_{i+1}^n)|$

$$|a^{\text{Roe}}(u_i^n, u_{i+1}^n)| = \sum_{\alpha=1}^{mP} |\lambda_\alpha| l_\alpha \otimes r_\alpha.$$

$$a^{\text{Roe}}(u_i^n, u_{i+1}^n) = \text{diag}([a^{\text{Roe}}(u_i^n, u_{i+1}^n)]^{SE}).$$

$$|a^{\text{Roe}}(u_i^n, u_{i+1}^n)| = \text{diag}(|[a^{\text{Roe}}(u_i^n, u_{i+1}^n)]^{SE}|).$$

⇒ Procedure **on each SE**.

Efficient approximation of $|a^{\text{Roe}}(u_i^n, u_{i+1}^n)|$

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⇒ Procedure **on each SE**.

1) Evaluate approximate eigenvalues of $[a^{\text{Roe}}(u_i^n, u_{i+1}^n)]^{SE}$

$$\{\lambda'_\gamma\} = \{\Lambda^k(U_{i,i+1}^{\text{Roe}}; \xi_\eta)\}.$$

Efficient approximation of $|a^{\text{Roe}}(u_i^n, u_{i+1}^n)|$

$$|a^{\text{Roe}}(u_i^n, u_{i+1}^n)| = \sum_{\alpha=1}^{mP} |\lambda_{\alpha}| l_{\alpha} \otimes r_{\alpha}.$$

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⇒ Procedure **on each SE**.

- 1) Evaluate **approximate eigenvalues** of $[a^{\text{Roe}}(u_i^n, u_{i+1}^n)]^{\text{SE}}$

$$\{\lambda'_{\gamma}\} = \{\Lambda^k(U_{i,i+1}^{\text{Roe}}; \xi_{\eta})\}.$$

- 2) Determine a local polynomial $q_{d, \{\lambda'_{\gamma}\}}^{\text{SE}}$ of low degree d ($d < 10$ in practice) **minimizing the least-square error**

$$\sum_{\gamma} (|\lambda'_{\gamma}| - q_{d, \{\lambda'_{\gamma}\}}^{\text{SE}}(\lambda'_{\gamma}))^2.$$

Efficient approximation of $|a^{\text{Roe}}(u_i^n, u_{i+1}^n)|$

$$|a^{\text{Roe}}(u_i^n, u_{i+1}^n)| = \sum_{\alpha=1}^{mP} |\lambda_\alpha| l_\alpha \otimes r_\alpha.$$

$$a^{\text{Roe}}(u_i^n, u_{i+1}^n) = \text{diag}([a^{\text{Roe}}(u_i^n, u_{i+1}^n)]^{SE}).$$

$$|a^{\text{Roe}}(u_i^n, u_{i+1}^n)| = \text{diag}(|[a^{\text{Roe}}(u_i^n, u_{i+1}^n)]^{SE}|).$$

⇒ Procedure **on each SE**.

- 1) Evaluate approximate eigenvalues of $[a^{\text{Roe}}(u_i^n, u_{i+1}^n)]^{SE}$

$$\{\lambda'_\gamma\} = \{\Lambda^k(U_{i,i+1}^{\text{Roe}}; \xi_\eta)\}.$$

- 2) Determine a local polynomial $q_{d,\{\lambda'_\gamma\}}^{SE}$ of low degree d ($d < 10$ in practice) minimizing the least-square error

$$\sum_\gamma (|\lambda'_\gamma| - q_{d,\{\lambda'_\gamma\}}^{SE}(\lambda'_\gamma))^2.$$

- 3) Approximation of $|[a^{\text{Roe}}(u_i^n, u_{i+1}^n)]^{SE}|$

$$|[a^{\text{Roe}}(u_i^n, u_{i+1}^n)]^{SE}| \approx q_{d,\{\lambda'_\gamma\}}^{SE}([a^{\text{Roe}}(u_i^n, u_{i+1}^n)]^{SE}).$$

The upwind scheme

$$u_i^{n+1} = u_i^n - \frac{\Delta^n t}{\Delta x} (\varphi(u_i^n, u_{i+1}^n) - \varphi(u_{i-1}^n, u_i^n)),$$

where the numerical flux $\varphi(u_i^n, u_{i+1}^n)$ is computed in this way

$$\varphi(u_i^n, u_{i+1}^n) = \frac{f(u_i^n) + f(u_{i+1}^n)}{2} - \text{diag}(q_d^{SE}([\mathbf{a}^{\text{Roe}}(u_i^n, u_{i+1}^n)])^{SE}) \frac{u_{i+1}^n - u_i^n}{2}.$$

Procedure on each SE.

CFL condition :

$$\frac{\Delta^n t}{\Delta x} = \min_{1 \leq SE \leq 2^{NNr}} \frac{CFL}{\max_{LR \in \mathcal{I}, \gamma} |\lambda'_{\gamma}|}.$$

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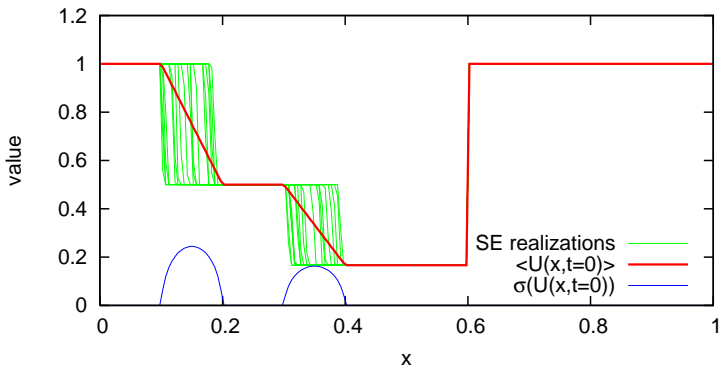
Entropy corrector

Periodic Burgers equation

$$\frac{\partial U}{\partial t} + \frac{\partial F(U)}{\partial x} = 0, \quad F(U) = \frac{U^2}{2},$$

initial random shock locations :

$$X_{1,2} = 0.1 + 0.1\xi_1, \quad X_{2,3} = 0.3 + 0.1\xi_2, \quad \xi_1, \xi_2 \sim \mathcal{U}[0, 1].$$

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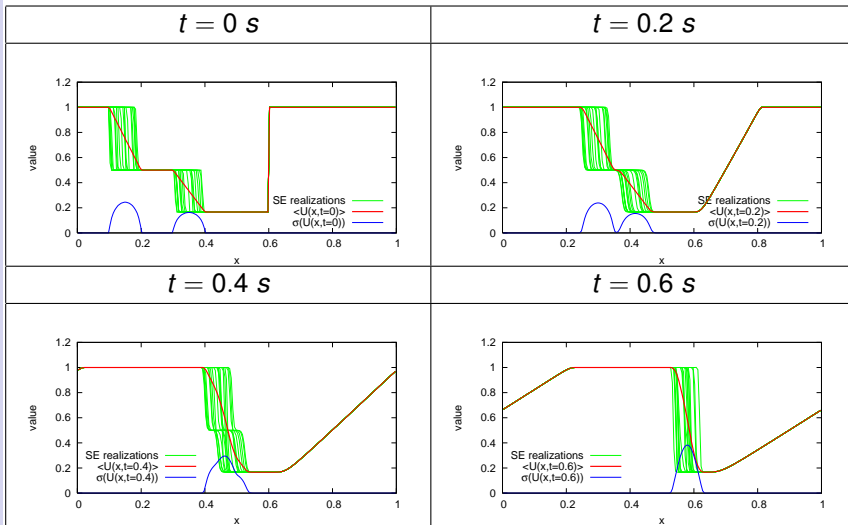
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Computations with $N_o = N_r = 3$.

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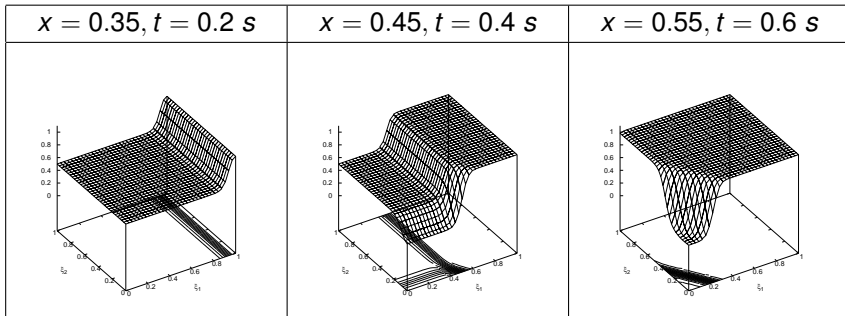
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Stochastic solution $U(x, t, (\xi_1, \xi_2))$ as a function of (ξ_1, ξ_2) for different spatial positions and times. $N_0 = N_r = 3$.

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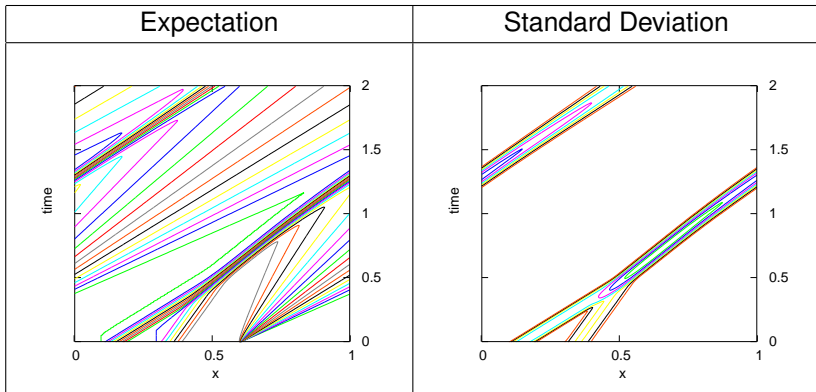
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Space-time isolines of the expectation $\langle U(x, t, \cdot) \rangle$ and standard deviation $\sigma(U(x, t, \cdot))$ of the stochastic solution. $N_0 = N_r = 3$.

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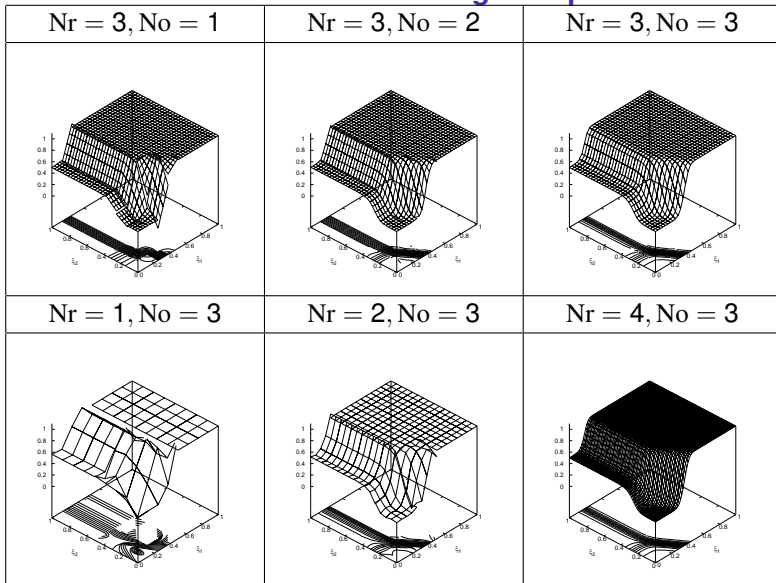
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Stochastic solution of the Burgers equation as a function of (ξ_1, ξ_2) at $x = 0.5$ and $t = 0.5$ for different N_r and N_o .

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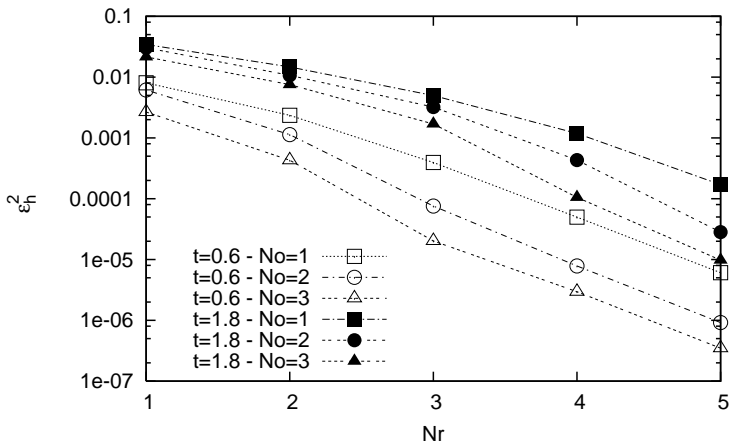
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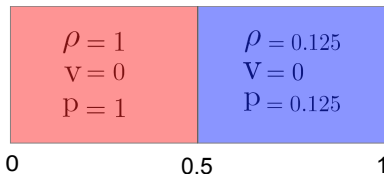
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$$\epsilon_h^2(t) := \frac{1}{M} \sum_{i=1}^M \int_{\Omega} \left(U_h^{\text{No}, \text{Nr}}(x, t, \xi^{(i)}) - U_h^{\text{MC}}(x, t, \xi^{(i)}) \right)^2.$$

Euler equations (Sod Shock Tube)



$$\frac{\partial U}{\partial t} + \frac{\partial F(U)}{\partial x} = 0,$$

$$U = (\rho, q, E)^T, \quad F(U) = (\rho v, \rho v^2 + p, v(E + p))^T,$$

$$v = \frac{q}{\rho}, \quad p = (\gamma - 1) \left(E - \frac{1}{2} \rho v^2 \right).$$

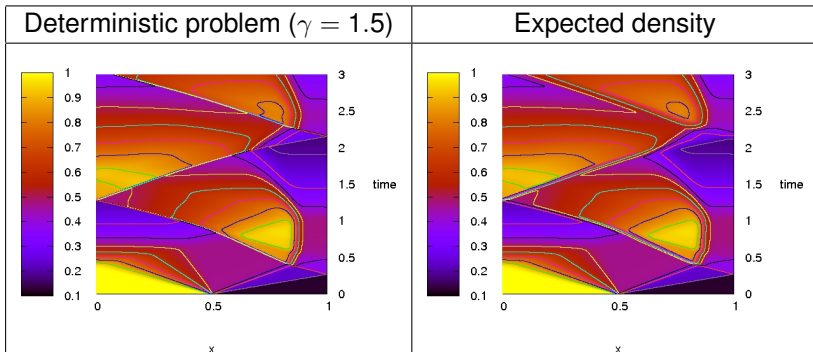
$$\gamma(\xi) = 1.4 + 0.2 \xi, \quad \xi \sim U[0, 1].$$

Euler equations

Computation of the Galerkin flux and Jacobian matrix : using a pseudo-spectral approximation (Debusschere et al. 04)

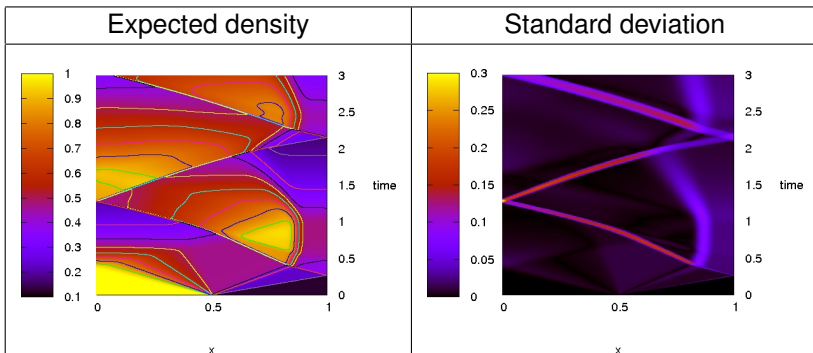
- $a \times b \approx a * b = \sum_{\alpha=0}^P (a * b)_{\alpha} \Psi_{\alpha}$,
 $(a * b)_{\alpha} = \sum_{\beta, \delta=0}^P a_{\beta} b_{\delta} \mathcal{M}_{\alpha\beta\gamma}$, $\mathcal{M}_{\alpha\beta\gamma} = \langle \Psi_{\alpha} \Psi_{\beta} \Psi_{\gamma} \rangle$
- $1/a \approx a^{-*}$ obtained by solving $a * a^{-*} = 1$
- $\sqrt{a} \approx a^{*/2}$ obtained by solving $(a^{*/2}) * (a^{*/2}) = a$
 $\rightarrow p \approx p^* = (\gamma - 1) * (E - ((q * q) * \rho^{-*})/2)$

Euler equations



Computations with $N_r = 3$ and $N_o = 2$.

Euler equations



Computations with $N_r = 3$ and $N_o = 2$.

Euler equations

Galerkin projection

Stochastic hyperbolic systems

Stochastic discretization

The Galerkin system

Hyperbolicity of the Galerkin system

Numerical method

Numerical scheme

Roe linearized matrix

Absolute value of the Roe matrix

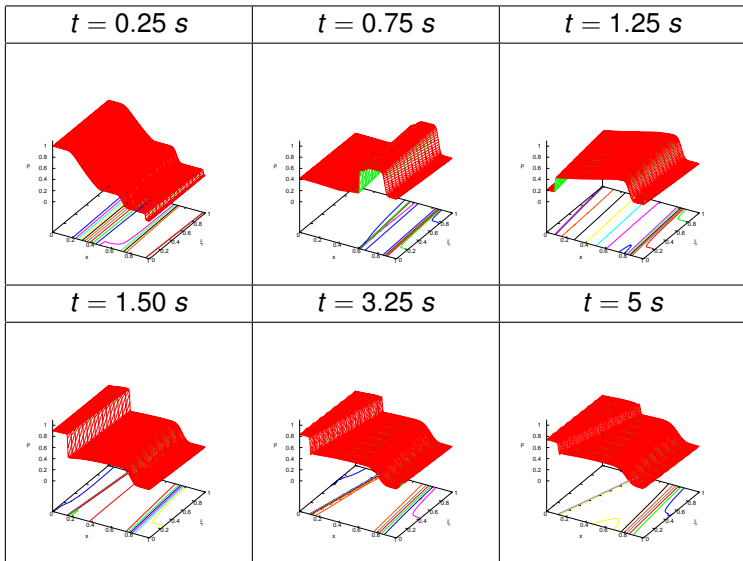
The upwind scheme

Results

Periodic Burgers equation

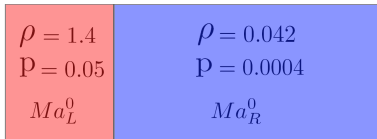
Euler equations

Entropy corrector



Stochastic density as a function of (x, ξ) . $N_r = 3$ and $N_o = 2$.

Entropy corrector Euler equations with sonic points



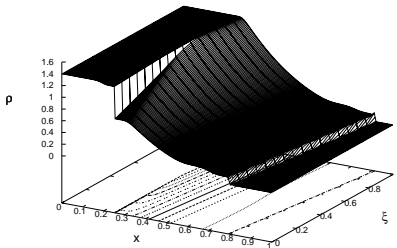
0

0.25

1

$$Ma^0(x, \xi) = \begin{cases} Ma_L^0(\xi), & x \in [0, 1/4], \\ Ma_R^0(\xi), & x \in]1/4, 1], \end{cases} \quad \xi \sim U[0, 1].$$

s.t. \exists sonic points for $\xi \in [0, 0.6]$.



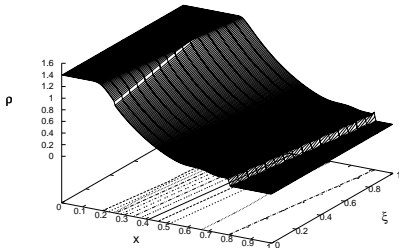
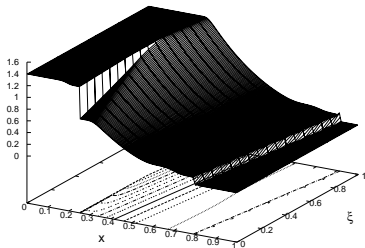
Entropy corrector

Non-parametrized entropy corrector proposed by Dubois and Mehlmann (96) for Roe solver in the deterministic case

- Nonlinear modification of the numerical flux in the vicinity of sonic points
- Detection of sonic expansion waves based on reconstruction of intermediate states for each couple of left and right states and test on sign of eigenvalues of the Roe linearized matrix

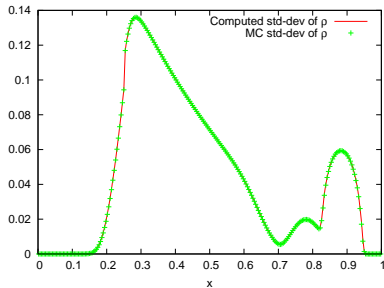
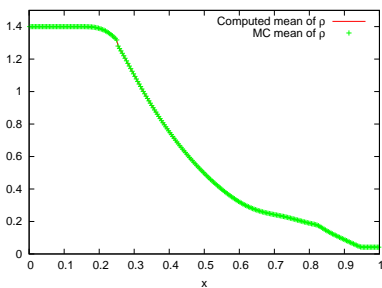
→ Adaptation to the present context using the approximate eigenvalues and eigenvectors of $a^{\text{Roe}}(u_i^n, u_{i+1}^n)$.

Entropy corrector Euler equations with sonic points



Stochastic density $\rho(x, t, \xi)$ at $t = 1$ obtained without (left) and with (right) the entropy corrector using $N_r = 3$ and $N_o = 2$.

Entropy corrector Euler equations with sonic points



Comparison of the mean and standard deviation of the numerical density at $t = 1$, computed with a Galerkin method (using $N_r = 3$ and $N_o = 2$) and a MC method.

Entropy corrector Euler equations with sonic points

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CPU improvements

- Only the eigenvalue $v - c$ can change its sign.
- Mean value averaged criterium \rightarrow for each interface LR , only test the entropy correction for the SE s.t.

$$E^{SE}[(v_L - c_L)] - ctol < 0 < E^{SE}[(v_R - c_R)] + ctol.$$

dim $\mathcal{S}^{Nr, No}$	No = 1, Nr = 3 16		No = 2, Nr = 3 24		No = 3, Nr = 3 32	
	T_{CPU}	R_{cor}	T_{CPU}	R_{cor}	T_{CPU}	R_{cor}
$ctol = +\infty$	11.7	49%	16.1	42%	21.6	38%
$ctol = 1e^{-1}$	8.2	19%	11.8	16%	16.4	14%
$ctol = 1e^{-2}$	6.5	4%	9.8	3%	13.9	2%
$ctol = 1e^{-3}$	6.1	<1%	9.3	<1%	13.5	<1%
ϵ_h	1.32e-3		7.17e-4		2.88e-4	

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Conclusion

- **intrusive stochastic Galerkin** method
- **Roe-type solver** with **upwind** matrices efficiently computed by an **original and fast method**
- **entropy correction** in the presence of sonic points only requiring marginal costs
- **accurate, stable** and **robust** method
- computational costs **scale as $\dim S^{\text{No}, \text{Nr}}$**
(at least for moderate No)

Under investigation :

savings in computational costs for problems with higher stochastic dimensions

→ **adaptive stochastic mesh refinement**

Adaptive stochastic mesh refinement Periodic Burgers equation

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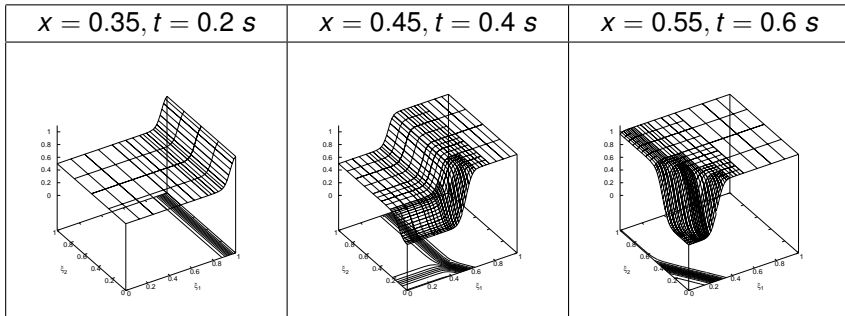
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Stochastic solution $U(x, t, (\xi_1, \xi_2))$ as a function of (ξ_1, ξ_2) for different spatial positions and times.

Adaptive stochastic mesh refinement Burgers equation

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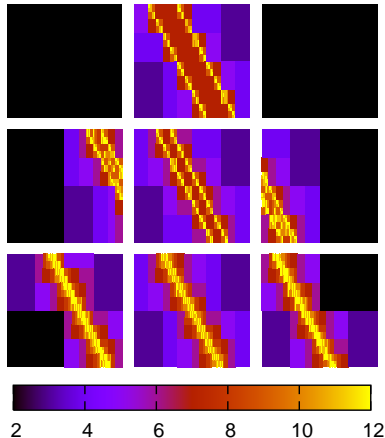
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Partition of the parameter space Ξ at times 0.2, 0.4 and 1.5 (from top to bottom) and spatial locations $x = 0.48, 0.5$ and 0.52 (from left to right). The color scale gives the depth of the leaves.

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References :

Intrusive Projection Methods with Upwinding for Uncertain
Nonlinear Hyperbolic Systems,

J. Tryoen, O. Le Maître, M. Ndjinga, A. Ern,
J. Comput. Phys., Under Review

Roe solver with Entropy Corrector for Uncertain Hyperbolic
Systems,

J. Tryoen, O. Le Maître, M. Ndjinga, A. Ern,
J. Comput. Appl. Math., Under Review

Euler equations

	Nr = 2		Nr = 3		Nr = 4	
	T_{CPU}	$\dim S^{Nr, No}$	T_{CPU}	$\dim S^{Nr, No}$	T_{CPU}	$\dim S^{Nr, No}$
No = 0	4.0	(4)	8.1	(8)	16.1	(16)
No = 1	6.9	(8)	13.9	(16)	27.8	(32)
No = 2	11.8	(12)	23.2	(24)	46.5	(48)
No = 3	17.1	(16)	34.1	(32)	68.1	(64)
No = 4	24.8	(20)	49.3	(40)	98.0	(80)

Normalized computational times T_{CPU} for different stochastic discretization parameters Nr and No. $N_c = 250$.

→ computational costs **scale as** $\dim S^{No, Nr}$ at least for moderate No.

Some details

- **Parallelisation** of the procedure on each **stochastic element** α_σ , $1 \leq \alpha_\sigma \leq P_\sigma$
- Compute the mP_π **characteristic variables** $\{\beta'_\gamma\}_{\gamma=1, \dots, mP_\pi}$

$$u_L - u_R \approx \sum_{\gamma=1}^{mP_\pi} \beta'_\gamma r'_\gamma(u_{LR}^{\text{Roe}}).$$

- Reconstruct the mP_π **intermediate states** at each physical interface

$$u'_\gamma = u'_{\gamma-1} + \beta'_\gamma r'_\gamma(u_{LR}^{\text{Roe}}).$$

- Determine the **set of sonic indices** :

$$S' = \{\gamma, \lambda'_\gamma(u'_{\gamma-1}) < 0 < \lambda'_\gamma(u'_\gamma)\}.$$

- The **indexing** of $\{\lambda'_\gamma\}_\gamma$ and $\{r'_\gamma\}_\gamma$, $\gamma = 1, \dots, mP_\pi$, provides a **correspondence** between approximate eigenvalues and eigenvectors and is **central** to determine S' .

Case of random variables with non-uniform distribution functions

Stochastic parametrization

$\xi = (\xi_1, \dots, \xi_N)$ vector of random variables with known
independent distribution functions.

Change of variables

$x(\xi) = (x_1(\xi_1), \dots, x_N(\xi_N)) = (p_1(\xi_1), \dots, p_N(\xi_N))$ with
 $(p_d(\xi_d))_{d=1, \dots, N}$ **cumulative density functions**

$$\rightarrow x(\xi) \sim \mathcal{U}([0, 1]^N).$$

Expansion of a process

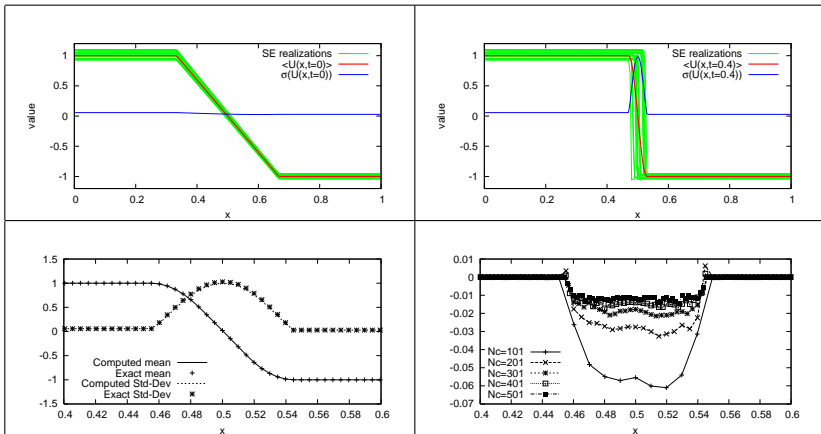
$$H(\xi) = \tilde{H}(x(\xi)) = \sum_{\alpha=1}^P \tilde{H}_{\alpha} \psi_{\alpha}(x(\xi)).$$

Burgers equation

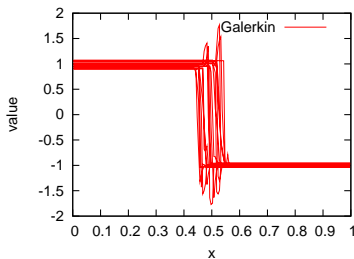
$$\frac{\partial U}{\partial t} + \frac{\partial F(U)}{\partial x} = 0, \quad F(U) = \frac{U^2}{2},$$

$$U^+(\xi_1) = 1 + 0.1(2\xi_1 - 1), \quad \xi_1 \sim \mathcal{U}[0, 1],$$

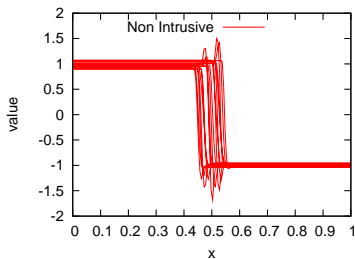
$$U^-(\xi_2) = -1 + 0.05(2\xi_2 - 1), \quad \xi_2 \sim \mathcal{U}[0, 1].$$



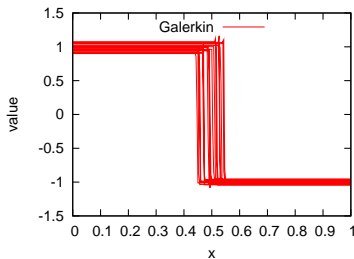
$N_o = 1$ and $N_r = 2$



$N_o = 1$ and $N_r = 2$



$N_o = 2$ and $N_r = 5$



$N_o = 2$ and $N_r = 5$

